Symmetry factorization of Selberg zeta functions and distribution of resonances

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#### Resonances for hyperbolic surfaces

 $(X, g) = \Gamma \setminus \mathbb{H}^2$  = geometrically finite hyperbolic surface of infinite area.

$$\mathcal{R}_X := \Big\{ \text{ poles of } R_X(s) := (\Delta_X - s(1-s))^{-1}, \text{ counted with multiplicity} \Big\}.$$

The first resonance occurs at  $s = \delta := \dim \Lambda(\Gamma)$ . [Pa

[Patterson: 1976, 1989]



There is a gap  $\varepsilon > 0$  such that  $\delta$  is the only resonance with Re  $s > \delta - \varepsilon$ . [Naud: 2005]

#### Resonance distribution conjectures

1. Fractal Weyl law [Sjostrand 1990, Lu-Sridhar-Zworski 2003]  $\# \left\{ \zeta \in \mathcal{R}_X : \operatorname{Re} \zeta \ge \sigma, |\operatorname{Im} \zeta - T| \le 1 \right\} \asymp T^{\delta} \quad \text{for some } \sigma < \delta.$ 

2. Essential spectral gap [Jakobson-Naud 2012]  $\inf \left\{ \sigma : \mathcal{R}_X \cap \{ \operatorname{Re} s > \sigma \} \text{ is finite} \right\} = \frac{\delta}{2}.$ 

3. Concentration of decay rates (quantum vs. classical)

$$\operatorname{Re}(\mathcal{R}_X)$$
 concentrates at  $\frac{\delta}{2}$ .

# Computations for 3-funnel surfaces

 $X(\ell_1, \ell_2, \ell_3)$  = hyperbolic pair of pants + funnel ends.



Strong analogies to 3-disk scattering systems



#### Resonance plots for X(10, 10, 10)



#### Resonance plots for X(10, 10, 10)





#### **Resonance chains**



Similar phenomena have been observed for symmetric *n*-disk scattering systems and open billiards, both experimentally and numerically.

# **Conjecture** Resonance chains are associated to (approximate) clustering of $\mathcal{L}_X$ on $\ell \cdot \mathbb{N}$ . [Barkhofen, Faure, Weich 2014].

Clustering for *X*(12, 12, 12):



## Fractal Weyl law



 $\# \{ \zeta \in \mathcal{R}_X : \operatorname{Re} \zeta \ge 0, \ 0 \le \operatorname{Im} \zeta \le t \} \text{ for } X(12, 14, 15)$ 

For  $\sigma > \delta/2$ ,

[Naud 2014]

$$\#\big\{\zeta\in\mathcal{R}_g: \operatorname{Re}\zeta\geq\sigma, \ 0\leq\operatorname{Im}\zeta\leq t\big\}=O(t^{1+\tau(\sigma)}),$$

with  $\tau(\sigma) < \delta$ .



# Spectral gap

#### **Envelope function**



#### Concentration of decay rates

Classical escape rate =  $1 - \delta$ . Quantum decay rate  $= \frac{1}{2} - \text{Re } \zeta$  (for a resonance at  $s = \zeta$ ).

 $\implies$  expect a concentration Re  $\mathcal{R}_X$  at  $\frac{\delta}{2}$ .



Histograms of Re  $\zeta$  for 0 Im  $\zeta$  < 20000.

# Selberg zeta function

For a geometrically finite surface of infinite area,  $X = \Gamma \setminus \mathbb{H}$ , the Selberg zeta function,

$$Z_X(s):=\prod_{\ell\in\mathcal{L}(\Gamma)}\prod_{k=1}^\inftyig(1-e^{-(s+k)\ell}ig),$$

converges for Re  $s > \delta$  and continues meromoprhically. [Patterson 1989, Guillopé 1992]

Divisor of 
$$Z_X(s) = \begin{cases} \text{zeros at } \mathcal{R}_X \\ \text{topological zeros at } -\mathbb{N}_0 & (\text{multiplicity } \propto \chi(X)) \\ \text{topological poles at } \frac{1}{2} -\mathbb{N}_0 & (\text{multiplicity } \propto n_{\text{cusps}}) \\ & [Patterson-Perry 2001, B-Judge-Perry 2005] \end{cases}$$

No cusps ( $\Gamma$  convex cocompact)  $\Longrightarrow$   $Z_X(s)$  is entire.

# Zeta function for Schottky groups

[Button 1998]:  $\Gamma$  convex co-compact  $\leftrightarrow$  classical Schottky group.



 $\Gamma$  is generated by  $S_1, \ldots, S_r \in PSL(2, \mathbb{R})$ , where

 $S_j$ : interior $(D_j) \rightarrow \text{exterior}(D_{j+r})$ .

#### Dynamical zeta function

Bowen-Series map: On  $U := \bigcup D_j$ , define

$$B|_{D_j} := S_j|_{D_j}$$

Ruelle transfer operator L(s) on  $L^2_{hol}(U)$ :

$$(L(s)u)(z) := \sum_{w \in U: Bw=z} B'(w)^{-s}u(w).$$

This setup is called a (holomorphic) iterated function scheme (IFS).

For  $\Gamma$  convex co-compact,

$$Z_X(s) = \det(1 - L(s)).$$

[Pollicott 1991]

#### Computation of the zeta function

The dynamical realization of the zeta function leads to a "cycle" expansion

$$Z_X(s) = 1 + \sum_{n=1}^{\infty} d_n(s),$$

with an estimate

$$|d_n(s)| \leq C e^{-c_1 n^2 - c_2 n \operatorname{Re} s + c_3 n |\operatorname{Im} s|}$$

[Cvitanović-Eckhardt 1989, Jenkinson-Pollicott 2002]

Here  $d_n(s)$  defined as a recursive sum over the functions

$$a_n(s) := -\frac{1}{n} \sum_{\sigma \in \mathcal{W}_n} \frac{e^{-s\ell(T_{\sigma})}}{1 - e^{-\ell(T_{\sigma})}},$$

where

$$\mathcal{W}_n := \Big\{ \sigma \in (\mathbb{Z}/2r\mathbb{Z})^n : \ \sigma_{j+1} \neq \sigma_j + r \text{ for } j = 1, \dots, n-1, \text{ and } \sigma_1 \neq \sigma_n + r \Big\},$$
  
and  $T_{\sigma} := S_{\sigma_1} \cdots S_{\sigma_n}.$ 

# Symmetry factorization

**Theorem** [B-Weich 2014] If a holomorphic IFS admits a finite symmetry group *G*, then the dynamical zeta function factors as a product over  $\hat{G}$ , the set of irreducible unitary representations of *G*.

(Complication: the natural action of G on  $L^2_{hol}(U)$  is not unitary.)

**Corollary** If a hyperbolic surface has finite symmetry group *G*, then

$$Z_X(s) = \prod_{\chi \in \widehat{G}} Z_X^{\chi}(s).$$

(Complication: the standard Schottky IFS may not reflect the full symmetry group of the surface.)

Symmetry factorization  $\rightsquigarrow$  very effective computation of  $Z_{\chi}^{\chi}$ .



Relative error terms at s = x + 1000i:

red = w/o symmetry reduction (using the first 170000 lengths) blue = symmetry reduced (using the first 47 lengths)

#### Resonances according to representation

 $X(7,7,7) \rightsquigarrow D_3 \times \mathbb{Z}_2$  symmetry



*I*<sub>1</sub>(dk blue), *I*<sub>2</sub>(lt blue), *II*<sub>1</sub>(red), *II*<sub>2</sub>(orange), *III*<sub>1</sub>(dk green), *III*<sub>2</sub>(lt green)

# Symmetric four-funnel surfaces

 $X(13, 13, 13, 13) \rightsquigarrow D_4 \times \mathbb{Z}_2$ 



 $I_1$ (dk blue),  $I_2$ (lt blue),  $II_1$ (red),  $II_2$ (orange),  $V_1$ (dark green),  $V_2$ (lt green),  $III_1$ ,  $III_2$ ,  $IV_1$ ,  $IV_2$  (none).

# Spectral gap

X(9,9,9)





# D<sub>2</sub> symmetry

X(6, 7, 7)



# Funneled torus





