Section 4.4: Vibrations. Harmonic oscillators

Main Topics:

- Mechanical vibrations (systems spring-mass)
- Harmonic oscillators
- Examples of second order differential equations.

Recall the equation describing the dynamics (or the vibrations) of a spring-mass system:

$$my'' + \gamma y' + ky = F$$

where

- the unknown function y = y(t) describes the motion of the mass at time t,
- *m* is the **mass**,
- k the spring constant,
- γ the damping factor,
- F = F(t) an **external force** applied on the system.

(*m*, *k* and γ are positive constants).

The initial conditions:

 $y(0) = y_0$ and $y'(0) = v_0$

(specifying the initial position y_0 and the initial velocity v_0) uniquely determine the motion.

- If F = 0 we say that the system is **unforced** or **free**.
- Free spring-mass systems are known as harmonic oscillators.

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Undamped free vibrations

This case corresponds to $\gamma = 0$ and F = 0.

The equation of the motion is

$$my'' + ky = 0$$

Setting

$$\omega_0^2 = k/m$$

this equation becomes

 $\mathbf{y}'' + \omega_0^2 \mathbf{y} = \mathbf{0} \,.$

Characteristic equation: $\lambda^2 + \omega_0^2 = 0$, with roots $\lambda = \pm i\omega_0$. The general solution is:

$$y(t) = A\cos(\omega_0 t) + B\sin(\omega_0 t)$$

for arbitrary constants A and B.

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The **phase-amplitude form** of the general solution $y(t) = A\cos(\omega_0 t) + B\sin(\omega_0 t)$ is

 $y(t) = R\cos(\omega_0 t - \delta)$



The graph of $y(t) = R \cos(\omega_0 t - \delta)$ is a displaced cosine function. This is a **periodic motion** (or **simple harmonic motion**) of the mass:

•
$$T = \frac{2\pi}{\omega_0}$$
 is the **period**,

- $\omega_0 = \sqrt{k/m}$ is the natural frequency,
- δ is the **phase**,
- R is the amplitude.

Example: Determine ω_0 , R and δ such that $y = -\cos t + \sqrt{3}\sin t$ can be written as $y = R\cos(\omega_0 t - \delta)$.

Damped free vibrations

This case corresponds to $\gamma \neq 0$ and F = 0.

The equation of the motion is

$$my'' + \gamma y' + ky = 0$$

Characteristic equation: $m\lambda^2 + \gamma\lambda + k = 0$ with roots:

$$\lambda_{1}, \lambda_{2} = \frac{-\gamma \pm \sqrt{\gamma^{2} - 4km}}{2m} = \frac{\gamma}{2m} \Big(-1 \pm \sqrt{1 - \frac{4km}{\gamma^{2}}} \Big)$$

The solutions of the equation of the motion are of three different types, depending on the sign of $\gamma^2 - 4km$.

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$$\lambda_{1}, \lambda_{2} = \frac{-\gamma \pm \sqrt{\gamma^{2} - 4km}}{2m} = \frac{\gamma}{2m} \Big(-1 \pm \sqrt{1 - \frac{4km}{\gamma^{2}}} \Big)$$

- Underdamped harmonic motion: $\gamma^2 4km < 0$ Complex conjugate roots $\lambda_1 = \mu + i\nu$, $\lambda_2 = \overline{\lambda_1}$. General solution: $y(t) = e^{-\gamma t/2m} (A\cos(\nu t) + B\sin(\nu t))$ with $\mu = -\gamma/2m < 0$ and $\nu = \frac{\sqrt{4km - \gamma^2}}{2m} > 0$.
- Critically damped harmonic motion: $\gamma^2 4km = 0$ Repeated *negative* roots $\lambda_1 = \lambda_2 = -\gamma/2m < 0$. General solution: $y(t) = (A + Bt)e^{-\gamma t/2m}$
- Overdamped harmonic motion: $\gamma^2 4km > 0$ Two distinct *negative* eigenvalues λ_1, λ_2 (because $\sqrt{\gamma^2 - 4km} < \gamma$). General solution: $y(t) = Ae^{\lambda_1 t} + Be^{\lambda_2 t}$

Remark: In all cases $\lim_{t\to+\infty} y(t) = 0$

(because of the damping, as t increases, the motion decreases and eventually stops)

A more detailed study of the underdamped harmonic motion: $\gamma^2 - 4km < 0$

General solution: $y(t) = e^{-\gamma t/2m} (A \cos(\nu t) + B \sin(\nu t))$ with $-\gamma/2m < 0$ and

$$\nu = \frac{\sqrt{4km - \gamma^2}}{2m} = \frac{\gamma}{2m} \sqrt{1 - \frac{4km}{\gamma^2}} > 0$$

As for the simple harmonic motion, set:

$$\begin{cases} R = \sqrt{A^2 + B^2} \\ A = R \cos \delta \\ B = R \sin \delta \end{cases}$$



Damped vibration; $y = Re^{-\gamma t/2m} \cos(\nu t - \delta)$. Figure 4.4.5 in J.Brannan & W. Joyce

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and write: $y(t) = Re^{-\gamma t/2m} \cos(\nu t - \delta)$

- the motion oscillates between the curves $y = Re^{-\gamma t/2m}$ (damped oscillation),
- ν is the quasi-frequency,
- $T_d = \frac{2\pi}{\nu}$ is the **quasi-period**.
- As $\gamma \to 2\sqrt{km}$ we have $\nu \to 0$ and $T_d \to \infty$ and there is no oscillation (critically damped motion). The value $\gamma = 2\sqrt{km}$ is called the **critical damping**.

Phase portraits for harmonic oscillators

Phase portraits are obtained from corresponding systems of first order differential equations:

$$\mathbf{x}' = \mathbf{A}\mathbf{x} = \begin{pmatrix} 0 & 1 \\ -k/m & -\gamma/m \end{pmatrix} \mathbf{x}$$

The origin is the unique equilibrium point because *A* is invertible (as det(*A*) = $k/m \neq 0$).

The roots of the characteristic equation $m\lambda^2 + \gamma\lambda + k = 0$ are the eigenvalues of **A**.

They give the nature of the equilibrium point of the different harmonic oscillators:

- a stable center for the undamped harmonic oscillator,
- a spiral sink for the underdamped harmonic oscillator,
- a **nodal sink** for both the overdamped and the critically damped harmonic oscillators.

Example:

A mass of weight 32 pounds is attached at the bottom end of a spring of natural lenght 6 ft. At the equilibrium, the spring measures 9.2 ft. The mass is initially released from rest at a point of 2 feet above the equilibrium position. Assume that the surrounding medium offers a resistance which is twice the instantaneous velocity.

- Write an initial value problem describing for the displacement *y* = *y*(*t*) of the mass as a function of the time *t*.
- Solve the initial value problem and determine *y*(*t*).
- Classify the type of harmonic oscillator. What are the parameters characterizing this motion?
- Describe the motion of the mass for very large values of *t*. Sketch the function *y*(*t*).
- Draw a phase portrait of the equivalent dynamical system. Give the type of equilibrium point.
- Describe the trajectory of the given initial value problem.